

MPS-based quantum impurity solvers

DMFT + DMRG

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Würzburg, 28 Feb 2015



Outline

Measure spectral functions

- Finite size \triangleright thermodynamic limit
- Analytically continue on real-time axis
- How to span the subspace using MPS?

DMFT with MPS

- Self-consistency – entanglement?
- Geometry of impurity problem – entanglement?

Results

- Equilibrium: two-patch DCA
- Non-equilibrium: quench from atomic limit

Why consider DMRG as impurity solver for DMFT?

Advantages over QMC

- EQ: direct access to frequency-dependent observables
- EQ/NEQ: no sign-problem

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Why hasn't it been used up to now?

- Lanczos: instable and imprecise [García, Hallberg & Rozenberg, PRL \(2004\)](#)
- DDMRG: extremely expensive [Nishimoto & Jeckelmann, JPhysCondMat, 2 papers \(2004\)](#),
[Karski, Raas & Uhrig, PRB \(2005\)](#), [Karski, Raas & Uhrig, PRB \(2008\)](#)
- Chebyshev and time evolution: much faster and precise [Ganahl, Thunström, Verstraete, Held & Evertz, PRB \(2014b\)](#), [Ganahl, Aichhorn, Thunström, Held, Evertz & Verstraete, arXiv \(2014a\)](#),
[Wolf, McCulloch, Parcollet & Schollwöck, PRB \(2014a\)](#), [Wolf, McCulloch & Schollwöck, PRB \(2014b\)](#)

Spectral functions: finite size \triangleright thermodynamic limit

Lin, Saad & Yang, arxiv:1308.5467 (2013)

e.g. Wolf, Justiniano, McCulloch & Schollwöck, arXiv:1501.07216 (2015)

Weiß, Wellein, Alvermann & Fehske, RMP 78, 275 (2006)

Continuous spectral function of thermodynamic limit from discrete spectral function of finite system?

$$A(\omega) = \langle \psi_0 | \delta(\omega - (H - E_0)) | \psi_0 \rangle, \quad |\psi_0\rangle = c^\dagger |E_0\rangle = \text{single-part. excit.}$$

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Expand $A(\omega)$ in “oscillatory functions” $p_n(\omega)$, truncate this expansion!

$$A(\omega) \simeq \sum_{n=0}^N c_n p_n(\omega) \quad \text{error} \simeq c_N \text{ for exp. conv.}$$

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- $p_n(\omega) = e^{i\omega \frac{n}{a}}$

$$\Rightarrow c_n = \langle \psi_0 | e^{-i(H-E_0)\frac{n}{a}} | \psi_0 \rangle \quad \text{time evolve}$$

- $p_n(\omega) = T_n(\omega)$

$$\Rightarrow c'_n = \langle \psi_0 | T_n(\frac{1}{a}(H - E_0)) | \psi_0 \rangle \quad \text{Chebyshev recurse}$$

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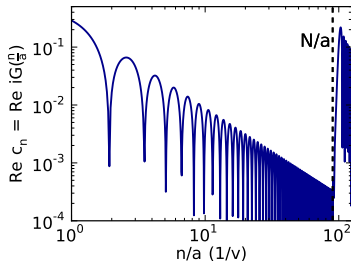
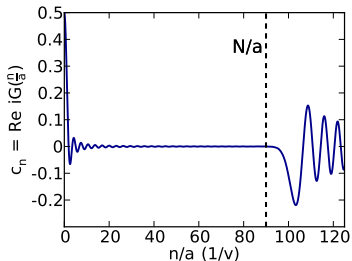
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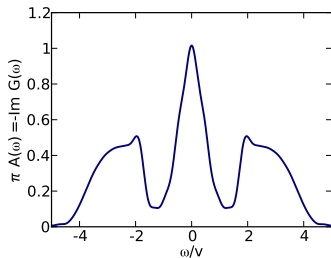
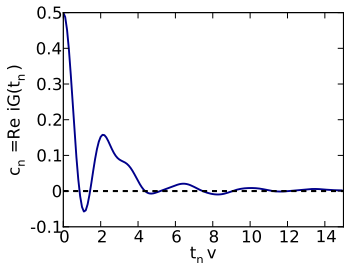
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Spectral functions: analytic continuation

Wolf, Justiniano, McCulloch & Schollwöck, arXiv:1501.07216 (2015)

Note that expansion coefficients are analytic functions of n for $n \in \mathbb{C}$

$$c_n = \langle \psi_0 | e^{-i(H-E_0)\frac{n}{a}} | \psi_0 \rangle \quad \text{exponential}$$

$$c'_n = \langle \psi_0 | \cos(n \arccos \frac{1}{a}(H - E_0)) | \psi_0 \rangle \quad \text{Chebyshev polynomial}$$

Find surrogate function $g(n)$ that agrees with c_n on $\{0, 1, \dots, N\}$

$\Rightarrow g(n)$ describes c_n also for $n \rightarrow \infty$.

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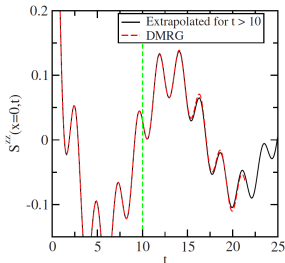
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Agreement in L_2 norm for short “times” ensures agreement for long times.

White & Affleck, PRB 77 134437 (2008)

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Fit function that accounts for the functional form of c_n

$$g(n) = \sum_j \alpha_j e^{i\omega_j n}$$

▷ *linear prediction*: recursive reformulation leads to *linear fit*

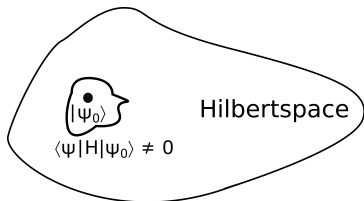
$$g(n) = \sum_j a_j g(n - j)$$

▷ for Chebyshev Ganahl, Thunström, Verstraete, Held & Evertz, Phys. Rev. B 90, 045144 (2014b)

Spectral functions: entanglement point of view

Wolf, Justiniano, McCulloch & Schollwöck, arXiv:1501.07216 (2015)

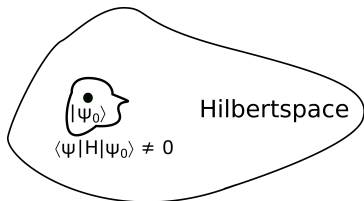
- Low entanglement of single-particle excitation $|\psi_0\rangle$
- Neighborhood of $|\psi_0\rangle$
 $\{|\psi\rangle : \langle\psi|H|\psi_0\rangle \neq 0\}$
- How to span this neighborhood?



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Immediate idea

- Orthogonal basis states “around” $|\psi_0\rangle$ most efficiently span this neighborhood (*Lanczos*)

$$|\psi_{n+1}\rangle = (H - \alpha_n)|\psi_n\rangle - \beta_n|\psi_{n-1}\rangle,$$

where α_n, β_n such that $\langle\psi_n|\psi_m\rangle = \delta_{nm}$

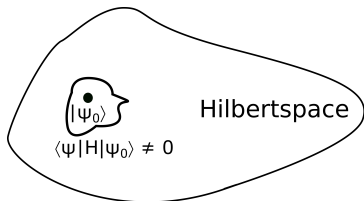
▷ But this simply is unstable, all the more, using MPS.

Dargel, Wöllert, Honecker, McCulloch, Schollwöck & Pruschke, PRB 85 205119 (2012)

Spectral functions: entanglement point of view

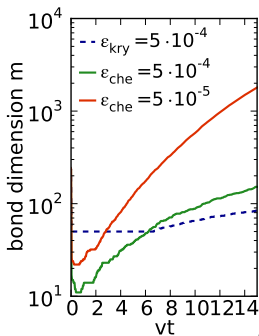
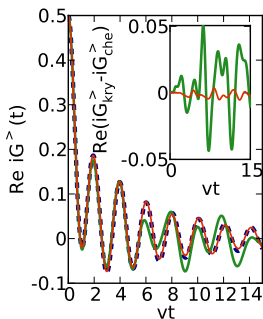
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Stable alternatives

- Chebyshev recurse from $|\psi_0\rangle$ \triangleright expand $A(\omega)$ in Chebyshev polynomials first MPS
[Holzner et al., PRB 83 195115](#)
- Time-propagate $|\psi_0\rangle$
 \triangleright Fourier expand $A(\omega)$

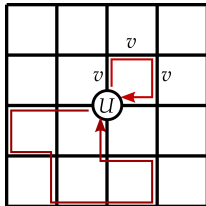


DMFT: self-consistency from entanglement point of view

Wolf, McCulloch & Schollwöck, PRB 90, 23513 (2014b)

Consider SIAM that is parametrized by the *hybridization function*

$$\Lambda(\omega) = v^2 G(\omega) \quad \text{or} \quad \Lambda(t) = v^2 G(t)$$

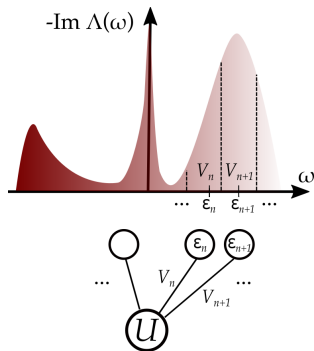


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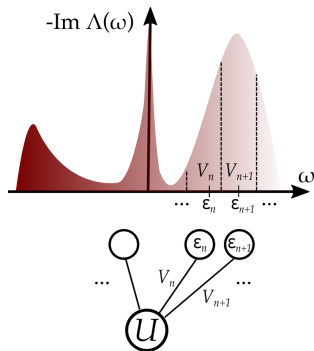
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▷ $|\psi(\infty)\rangle$ is highly entangled. DDMRG *a priori* involves this point.



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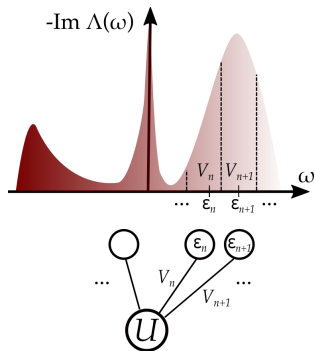
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Remark (i) If Fourier transform can be avoided: *much* easier!

▷ NEQDMFT starting from uncorrelated initial state solved on time-slices. Gramsch, Balzer, Eckstein & Kollar, PRB 88, 235106 (2013)

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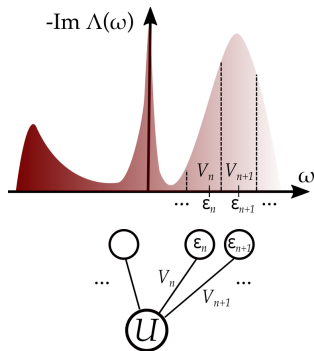
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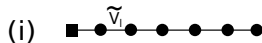
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Remark (ii) When solving the self-consistency iteratively in frequency space: successively increase resolution by going from short times to longer times ▷ exponential speed-up

DMFT: entanglement from geometry point of view

Wolf, McCulloch & Schollwöck, PRB 90, 23513 (2014b)

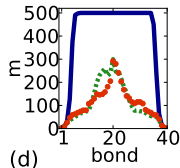
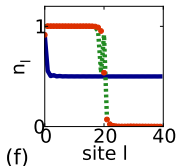
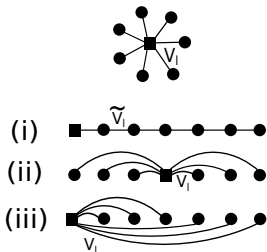
▷ Lower entanglement: star or chain geometry?



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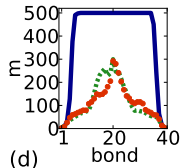
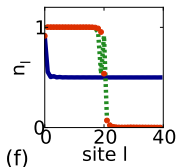
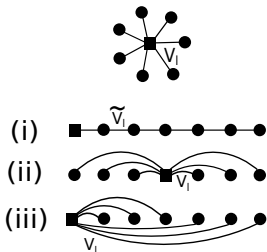
quickly converging DMRG algorithm

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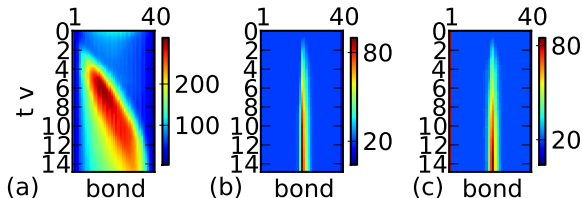
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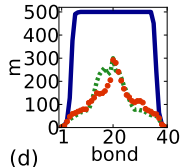
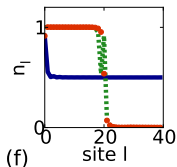
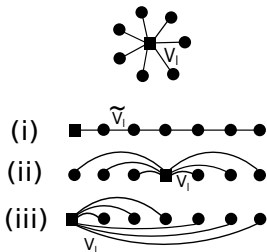
▷ Highly different bond dimension growth in different geometries



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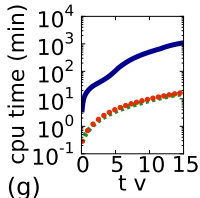
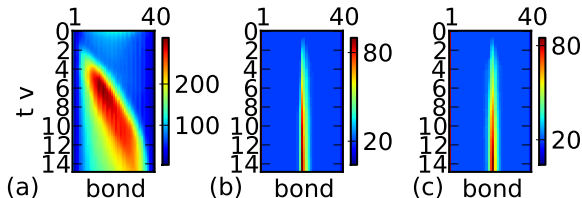
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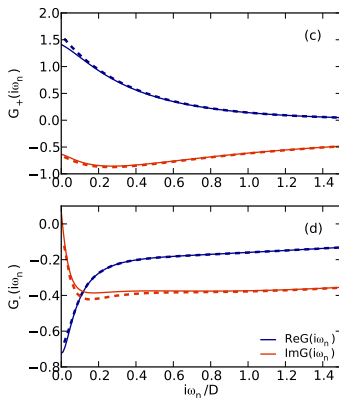
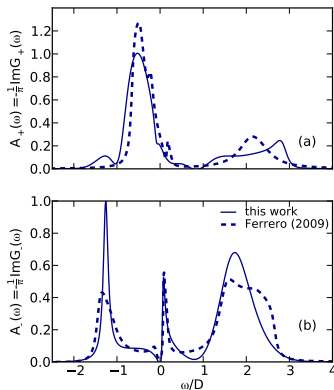


Results: two-site cluster DCA

Wolf, McCulloch, Parcollet & Schollwöck, PRB 90, 115124 (2014a)

CTQMC by Ferrero, Cornaglia, De Leo, Parcollet, Kotliar & Georges, PRB 80, 064501 (2009)

Model: Hole-doped Hubbard model on 2 dim. square lattice



▷ Pseudo-gap well reproduced

Results: non-equilibrium DMFT

MPS solution [Wolf, McCulloch & Schollwöck, PRB 90, 23513 \(2014b\)](#)

Hamiltonian representation [Gramsch, Balzer, Eckstein & Kollar, PRB 88, 235106 \(2013\)](#)

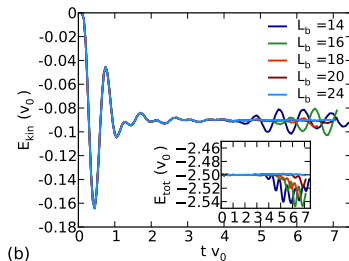
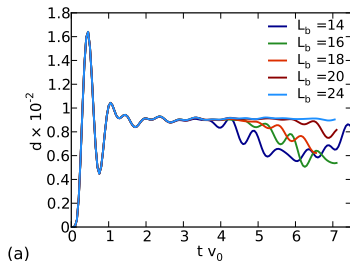
Model: single-band Hubbard model on Bethe lattice

▷ quench from atomic limit $v = 0$ to $v = v_0$

Strong interactions: $U = 10v_0$

Exact diagonalization: $t_{\max} \sim 3/v_0$

MPS: $t_{\max} \sim 7/v_0$



Results: non-equilibrium DMFT

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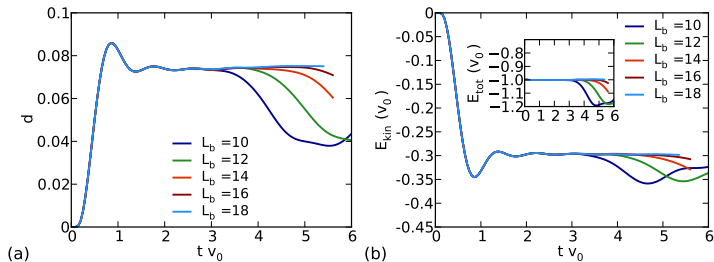
Model: single-band Hubbard model on Bethe lattice

▷ quench from atomic limit $v = 0$ to $v = v_0$

Intermediate interactions: $U = 4v_0$

Exact diagonalization: $t_{\max} \sim 3/v_0$

MPS: $t_{\max} \sim 5.5/v_0$



Summary and Outlook

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- Exploit the self-consistency from an entanglement point of view
- Star geometry much less entangled than chain geometry

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Thank you!

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